# Nonlinear connections and the problem of metrizability 

By J. SZILASI (Debrecen) and Z. MUZSNAY (Debrecen)

Dedicated to Professor Lajos Tamássy on his 70th birthday

## 1. Introduction

(1.1) A nonlinear connection in a vector bundle is nothing but a vector 1 -form, which is a projector whose kernel is the vertical bundle. This paper is devoted to a study of some problems concerning such connections. Sections 2 and 3 are of preparatory nature: we briefly summarize the necessary tools of Frölicher-Nijenhuis formalism of vector valued forms and collect some fundamental facts concerning differential operators and their formal integrability ("Spencer-Quillen-Goldschmidt technique"). In section 4 we present some basic facts about nonlinear connections. We show that these can be characterized as quasi-scalar (nonlinear) differential operators and that their tension is in fact a generalization of the deflection known from Matsumoto's theory of Finsler connections. The homogeneity condition for nonlinear connections (which results in linear connections) is also expressed in a new, illuminating way. In section 5 we formulate the problem of metrizability and show that the question is the formal integrability of a first order linear differential operator arising from the connection. In section 6 we show that our operator is involutive, while in the concluding section 7 we demonstrate that in the flat case all of the conditions of the modern version of the Cartan-Kähler theorem are satisfied so that the metrizability problem always has a solution. Of course, the vanishing of the curvature is a very restrictive assumption, but we hope that this application of the Spencer-Quillen-Goldschmidt technique has enough interest and (possibly) the method used here may also be successful in more general and much more complicated cases.
(1.2) Before we embark on the actual topic of the paper, we wish to review briefly the basic conventions and notations that will be used in the
sequel. - We shall always be working in the caterory of the finite dimensional, second countable, smooth manifolds. The tangent and cotangent bundle of a manifold $M$ will be written as a triple $\tau_{M}=(T M, \tau, M)$ and $\tau_{M}^{*}=\left(T^{*} M, \tau^{*}, M\right)$, resp.; $T$ is the tangent functor. $C^{\infty}(M)$ denotes the algebra of smooth functions on $M ; \mathfrak{X}(M)$ is the $C^{\infty}(M)$-module of vector fields on $M ; \wedge(M)$ and $S(M)$ are the exterior and the symmetric algebra over $M$. As usual, $d$ is the operator of the exterior derivative, while $i_{X}: \wedge(M) \rightarrow \wedge(M)(X \in \mathfrak{X}(M))$ is the substitution operator (or "interior product by $X "$ ).

In coordinate calculations Einstein's summation convention will be used as appropriate.

## 2. Frölicher-Nijenhuis formalism

Let $N$ be a manifold, $k \in \mathbb{N}$. A vector valued $k$-form (briefly a vector $k$-form) is a skew-symmetric $C^{\infty}(N)$-multilinear mapping $K: \mathfrak{X}^{k}(N) \rightarrow$ $\mathfrak{X}(N)$, with the agreement that the vector 0 -forms are the elements of $\mathfrak{X}(N)$. Frölicher-Nijenhuis theory associates to any vector $k$-form $K$ two derivations of the exterior algebra $\wedge(N)$ : a derivation of degree $k-1$ which is denoted by $i_{K}$ and a derivation of degree $k$, denoted by $d_{K} . i_{K}$ and $d_{K}$ reduce to the above substitution operator $i_{X}$ and the Lie derivative $d_{X}$ in case of vector 0 -forms (i.e. vector fields), while in general they can be described as follows:
(i) Derivations $i_{K}$ annihilate $C^{\infty}(N)$ and are determined by their actions on the (scalar) 1-forms:

$$
\begin{gather*}
\forall \omega \in \wedge^{1}(N), \quad X_{i} \in \mathfrak{X}(N) \quad(1 \leq i \leq k): \\
i_{K} \omega\left(X_{1}, \ldots, X_{k}\right)=\omega\left(K\left(X_{1}, \ldots, X_{k}\right)\right) . \tag{1}
\end{gather*}
$$

(ii) $d_{K}=\left[i_{K}, d\right]:=i_{K} \circ d+(-1)^{k} d \circ i_{K}$, so - in particular -

$$
\begin{equation*}
\forall f \in C^{\infty}(N): d_{K} f=i_{K} d f \tag{2}
\end{equation*}
$$

this last relation determines $K$ uniquely. (In (i) and (ii) $k:=\operatorname{deg} K$ is the degree of $K$.)

The bracket of two derivations $d_{K}, d_{L}$ is introduced thus:

$$
\begin{equation*}
\left[d_{K}, d_{L}\right]=d_{K} \circ d_{L}-(-1)^{k \ell} d_{L} \circ d_{K} \tag{3}
\end{equation*}
$$

( $k=\operatorname{deg} K, \ell=\operatorname{deg} L$ ). A fundamental result of the theory states that for each vector $k$-form $K$ and vector $\ell$-form $L$ there exists a unique $k+\ell$-form [ $K, L]$ such that

$$
d_{[K, L]}=\left[d_{K}, d_{L}\right] .
$$

[K,L] is called the Frölicher-Nijenhuis bracket, briefly the bracket of $K$ and $L$. In case of vector 0 -forms it reduces to the usual bracket of vector fields. In the general case, the evaluation of $[K, L]$ on vector fields is quite complicated, for an effective formula we refer to [10]. Fortunately, in this paper we shall have to evaluate brackets only in the following very simple situations:
(i) $K$ is a vector 1-form, $L:=Z \in \mathfrak{X}(N)$. Then

$$
\begin{equation*}
\forall X \in \mathfrak{X}(N):[K, Z](X)=[K(X), Z]-K[X, Z] . \tag{4}
\end{equation*}
$$

(ii) Let $K$ be a vector 1-form again. By (3) we get that

$$
\left[d_{K}, d_{K}\right]=2 d_{K} \circ d_{K}=d_{[K, K]}, \text { so }
$$

$$
\begin{equation*}
d_{K}^{2}=d_{\frac{1}{2}[K, K]} \tag{5}
\end{equation*}
$$

Here $N_{K}:=\frac{1}{2}[K, K]$ is called the Nijenhuis torsion of $K$. It can be characterized by the formula $N_{K}(X, Y)=[K(X), K(Y)]+$ $K \circ K[X, Y]-K[K(X), Y]-K[X, K(Y)] \quad(X, Y \in \mathfrak{X}(N))$.

## 3. Differential operators and formal integrability

(3.1) Suppose that $M$ is an $n$-dimensional manifold and consider a vector bundle $\xi=(E, p, M)$ of rank $r$ over $M$. Let Sec $\xi$ be the $C^{\infty}(M)$ module of sections of $\xi$ and $J E$ be the 1-jet space of $\xi$, i.e.

$$
J E=\left\{j_{x} \sigma \mid x \in M, \sigma \in \operatorname{Sec} \xi\right\}
$$

where $j_{x} \sigma$ is the 1 -jet of $\sigma$ at $x$. We denote by $\pi$ the target projection

$$
J E \rightarrow E, \quad j_{x} \sigma \mapsto \sigma(x)
$$

It is well-known that the sequence

$$
0 \rightarrow T^{*} M \otimes E \xrightarrow{\varepsilon} J E \xrightarrow{\pi} E \rightarrow 0,
$$

where

$$
\varepsilon: d f(x) \otimes \sigma(x) \longmapsto j_{x}[(f-f(x)) \sigma]
$$

$\left(f \in C^{\infty}(M), \sigma \in \operatorname{Sec} \xi, x \in M\right)$ is a short exact sequence of vector bundles. More generally (see e.g. [4], Ch. VI, $\S 1)$, if $J_{k} E$ is the space of $k$-jets $\left(k \in \mathbb{N}^{+}\right)$,

$$
\pi_{k}: J_{k} E \rightarrow J_{k-1} E, \quad j_{k} \sigma \longmapsto j_{k-1} \sigma
$$

is the natural projection ( $\pi_{0}:=\pi, j_{0} \sigma:=\sigma$ ) and

$$
\varepsilon: S^{k} T M \times E \rightarrow J_{k} E
$$

is defined by

$$
\varepsilon\left(d f_{1}(x) \odot \cdots \odot d f_{k}(x) \otimes \sigma(x)\right):=\left[j_{k}\left(f_{1} \ldots f_{k} \sigma\right)\right]_{x}
$$

$\left(f_{i} \in C^{\infty}(M), f_{i}(x)=0,1 \leq i \leq k ; \odot\right.$ is the symbol of symmetric product), then the sequence

$$
0 \rightarrow S^{k} T M \otimes E \rightarrow J_{k} E \xrightarrow{\varepsilon} J_{k-1} E \xrightarrow{\pi_{k}} 0
$$

is also a short exact sequence of vector bundles.
(3.2) Let $\eta=(F, \rho, M)$ be another vector bundle over $M$. - We recall that a mapping $P: \operatorname{Sec} \xi \rightarrow \operatorname{Sec} \eta, \sigma \mapsto P \sigma$ is said to be a linear differential operator of order $k\left(\in \mathbb{N}^{+}\right)$if

$$
\forall x \in M, \sigma \in \operatorname{Sec} \xi:\left(j_{k} \sigma\right)_{x}=0 \Longrightarrow P \sigma(x)=0
$$

Any differential operator of order $k$ can be identified with an $M$-morphism $\bar{P}: J_{k} E \rightarrow F$ such that $P=\bar{P}_{*} \circ j_{k}$, where $\bar{P}_{*}: \operatorname{Sec} J_{k} \xi \rightarrow \operatorname{Sec} \eta$ is induced by $\bar{P}$ according to [6], p.63. We shall use this observation without any comment.

If $\operatorname{rank} \bar{P}$ is constant and $R_{k}:=\operatorname{Ker} \bar{P}$, then the sequence

$$
0 \rightarrow R_{k} \xrightarrow{\text { inc }} J_{k} E \xrightarrow{\bar{P}} F \rightarrow 0
$$

is clearly exact (inc:= inclusion).
The $\ell$-th prolongation $p_{\ell}(P)$ of $P$ is defined as follows:

$$
p_{\ell}(P): J_{k+\ell} E \rightarrow J_{\ell} F, \quad j_{k+\ell}(\sigma) \mapsto j_{\ell}(P \sigma)
$$

Let $R_{k+\ell}:=\operatorname{Ker} p_{\ell}(P)$. The natural projection $\pi_{k+\ell}: J_{k+\ell} E \rightarrow J_{k} E$ induces a mapping $R_{k+\ell} \rightarrow R_{k}$ by restriction, it will be denoted also by $\pi_{k+\ell}$. - Keeping these notations, we can formulate the following fundamental

Definition 1. A $k$-th order differential operator is said to be formally integrable if $\forall \ell \in \mathbb{N}^{+}: \pi_{k+\ell}: R_{k+\ell} \rightarrow R_{k}$ is surjective.

Other key notions are explained in the next
Definition 2.
(i) The symbol of a differential operator $P: J_{k} E \rightarrow F$ is the mapping

$$
s_{0}(P):=P \circ \varepsilon: S^{k} T^{*} M \otimes E \rightarrow F .
$$

The symbol of the $\ell$-th prolongation operator $p_{\ell}(P): J_{k+\ell} E \rightarrow J_{\ell} F$ is the mapping $s_{\ell}(P): S^{k+\ell} T^{*} M \otimes E \rightarrow S^{\ell} T^{*} M \otimes F$ introduced by the diagram
(ii) Suppose, in particular, that $P: \operatorname{Sec} \xi \rightarrow \operatorname{Sec} \eta$ is a first-order linear differential operator. Then

$$
s_{0}(P): T_{x}^{*} M \otimes E_{x} \rightarrow E_{x}, \quad(d f)_{x} \otimes \sigma(x) \mapsto P(f \sigma)(x)
$$

(where $f(x)=0$ ). Fixing an element of $T_{x} M$, we get a partial mapping $E_{x} \rightarrow E_{x} .-P$ is said to be quasi-scalar if each of these mappings means multiplication by a scalar (cf. [12], def. 19.33).

The third important preparatory step is summarized in
Definition 3. Consider a linear differential operator $P: J_{k} E \rightarrow F$. Let $\forall \ell \in \mathbb{N}: g_{\ell}:=\operatorname{Ker} s_{\ell}(P)$. A basis $\left(e_{i}\right)_{1 \leq i \leq n}$ of $T_{x} M$ is said to be quasi-regular if

$$
\operatorname{dim}\left(g_{1}\right)_{x}=\operatorname{dim}\left(g_{0}\right)_{x}+\sum_{j=1}^{n-1} \operatorname{dim}\left(g_{0}\right)_{\left(e_{1}, \ldots, e_{j}\right)},
$$

where

$$
\left(g_{0}\right)_{\left(e_{1}, \ldots, e_{j}\right)}:=\left\{A \in g_{0} \mid i_{e_{1}} A=\cdots=i_{e_{j}} A=0\right\} \quad(1 \leq j \leq n-1)
$$

If for each $x \in M$ there exists a quasi-regular basis of $T_{x} M$ then $P$ is called involutive.

With the help of the terms introduced, the classical Cartan-Kähler theorem (see e.g. [3], Th. 18.13.8) can be translated into the following extremely compact and elegant form:

Cartan-Kähler-Goldschmidt theorem. Let $P: J_{k} E \rightarrow F$ be a $k$-th order linear differential operator. If
$1^{\circ} P$ is involutive, and
$2^{\circ} \pi_{k+1}: R_{k+1} \rightarrow R_{k}$ is surjective, then $P$ is formally integrable.
For a comprehensive, but well-readable treatment of the subject we refer to [1] and [13]. A very good brief survey is presented in [5]; see also [2].
(3.3) To conclude this section, we take a look at the nonlinear case. - A first order nonlinear differential operator, briefly differential operator from $\xi$ to $\eta$ is simply a mapping $P: \operatorname{Sec} \xi \rightarrow \operatorname{Sec} \eta$ which has a factorisation of the form $P=\bar{P}_{*} \circ j$, where $\bar{P}: J E \rightarrow F$ is just fibre-preserving. Taking Gâteaux-derivative, nonlinear differential operators can be linearized. Namely, if $P: \operatorname{Sec} \xi \rightarrow \operatorname{Sec} \eta$ is a nonlinear differential operator, $\sigma \in \operatorname{Sec} \xi$ is fixed and

$$
P_{\sigma}^{\prime}(\rho):=\left.\frac{d}{d t} P(\sigma+t \rho)\right|_{t=0}
$$

then $P_{\sigma}^{\prime}$ is a linear differential operator which is called the linearized operator of $P$ along $\sigma$.

We show a simple but very useful example. - Let $X \in \mathfrak{X}(E)$ be a projectable vector field:

$$
X \stackrel{P}{\sim} \bar{X}, \quad \bar{X} \in \mathfrak{X}(M) .
$$

Then the mapping

$$
P_{X}: \operatorname{Sec} \xi \rightarrow \operatorname{Sec} \xi, \sigma \mapsto \alpha \circ(T \sigma \circ \bar{X}-X \circ \sigma)
$$

$(\alpha: V E \rightarrow E$ is the well-known canonical surjection; see e.g. [6], p.291) is a nonlinear differential operator. Its linearized operator along a section $\sigma$ acts as follows:

$$
\forall \rho \in \operatorname{Sec} \xi:\left(P_{X}^{\prime}\right)_{\sigma}(\rho)=\alpha \circ\left[X, \rho^{v}\right] \circ \sigma,
$$

where $\rho^{v}:=\alpha^{\sharp} \rho$ is the vertical lift of $\rho$ (see [18], Lemma 3). As for the symbol of $P_{X}$, an easy calculation yields the formula

$$
\begin{gathered}
s_{0}\left(P_{X}\right)_{\sigma}^{\prime}(d f \otimes \rho)_{x}=(\bar{X} f)^{v} \sigma(x) \rho(x) \\
\left((\bar{X} f)^{v}=\bar{X} f \circ p\right),
\end{gathered}
$$

from which it follows at first sight that $\left(P_{x}\right)_{\sigma}^{\prime}$-and consequently $P_{X}-$ is quasi-scalar.

## 4. Nonlinear connections

(4.1) Let the above vector bundle $\xi=(E, p, M)$ be given. We shall write its vertical bundle as the triple $V \xi=\left(T^{v} E, p_{v}, E\right) ; \mathfrak{X}^{v} E:=\operatorname{Sec} V \xi$ is the module of vertical vector fields.

Definition 4. A nonlinear connection, briefly a connection in $\xi$ is a vector 1-form $h: \mathfrak{X}(E) \rightarrow \mathfrak{X}(E)$, satisfying the following conditions:
$1^{\circ} h^{2}=h$, i.e. $h$ is a projector;
$2^{\circ}$ Ker $h=\mathfrak{X}^{v} E$.
Of course, this is only a possible one among the numerous definitions of a connection, which seems preferable to us only for practical reasons. Let $\mathcal{H}(\xi)$ be the set of all connections in $\xi$. Then $\mathcal{H}(\xi)$ is an affine bundle ([14]), Def. 2.4.4) in a natural manner. Clearly, each $h \in \mathcal{H}(\xi)$ gives rise to direct decompositions $\mathfrak{X}(E)=\mathfrak{X}^{v} E \oplus \operatorname{Im} h$ (in the module sense) and $T E=T^{v} E \oplus H^{v} E$ (in the vector bundle sense), where $H \xi:=\left(T^{h} E, p_{h}, E\right)$ is the so-called horizontal subbundle belonging to $h$, such that

$$
\operatorname{Sec} H \xi=\operatorname{Im} h=: \mathfrak{X}^{h} E .
$$

Elements of $\mathfrak{X}^{h} E$ are the horizontal vector fields of the connection. Finally, it is easy to show that there exists a unique mapping

$$
\ell^{h}: \mathfrak{X}(M) \rightarrow \mathfrak{X}^{h} E, \quad X \mapsto X^{h}
$$

characterized by the condition $X^{h} \stackrel{p}{\sim} X$; it is called horizontal lifting.

## Proposition 1.

(a) Let $h \in \mathcal{H}(\xi), X \in \mathfrak{X}(M), \nabla_{X}:=P_{X^{h}}$ (see (3.3)). Then $\nabla_{X}$ is a quasi-scalar differential operator, said to be covariant derivative by $X$ with respect to $h$.
(b) Consider the bundle $V_{M} \xi:=\left(T^{v} E, p \circ p_{v}, M\right)$ and suppose that a mapping $X \in \mathfrak{X}(M) \mapsto P_{X}$ is given, such that
(i) $\forall X \in \mathfrak{X}(M): P_{X}: \operatorname{Sec} \xi \rightarrow \operatorname{Sec} V_{M} \xi$ is a quasi-scalar (nonlinear) differential operator,
(ii) $\forall \sigma \in \operatorname{Sec} \xi: p_{v} \circ P_{X} \sigma=\sigma$,
then there exists a unique connection $h \in \mathcal{H}(\xi)$ for which $\forall X \in \mathfrak{X}(M)$ : $\nabla_{X}=\alpha \circ P_{X}$.

Proof. (a) follows immediately by (3.3). We sketch, how one can get the converse statement (b). - Observe first that condition (i) is meaningful because $P_{X}$ can be identified with the operator $\alpha \circ P_{X}: \operatorname{Sec} \xi \rightarrow \operatorname{Sec} \xi$ canonically. To become conscious of this, let us define the mapping

$$
\ell^{h}: \mathfrak{X}(M) \rightarrow \mathfrak{X}(E), \quad X \mapsto \ell^{h}(X)=: X^{h}
$$

as follows:

$$
X^{h}(z):=T_{x}[X(x)]-\left(P_{X} \sigma\right)(x), \quad \text { if } z=\sigma(x), \sigma \in \operatorname{Sec} \xi
$$

Applying condition (i), one can check by a direct calculation that $X^{h}(z)$ does not depend on the section for which $z=\sigma(x)$, so $X^{h} \in \mathfrak{X}(E)$ is a well-defined vector field. Since $\left(P_{X} \sigma\right)(x) \in T_{\sigma(x)}^{V} E:=\operatorname{Ker} T_{\sigma(x)} P$, $X^{h} \stackrel{p}{\sim} X$. Finally, the introduced mapping $\ell^{h}$ determines uniquely a connection $h \in \mathcal{H}(\xi)$ such that $\operatorname{Im} \ell^{h}=\mathfrak{X}^{h} E$ and $\forall X \in \mathfrak{X}(M): \nabla_{X}=\alpha \circ P_{X}$.

Proposition 1 characterizes (nonlinear) connections as quasi-scalar (nonlinear) differential operators, extending the similar characterization of linear connections; cf. [11], $\S 9$.
(4.2) Let $C: E \rightarrow T^{v} E$ be the Liouville vector field on $E$.

Definition 5. (cf. [7]) The tension of a connection $h \in \mathcal{H}(\xi)$ is the vector 1 -form $T:=[h, C]$. In case of $T=0$ we say that $h$ satisfies the homogeneity condition (HC) and speak of a linear connection.

Note that by (4)

$$
\begin{aligned}
(H C) & \Longleftrightarrow \forall X \in \mathfrak{X}(E):[h X, C]-h[X, C]=0 \\
& \left.\Longleftrightarrow \forall X \in \mathfrak{X}(M):\left[X^{h}, C\right]=0 \quad \text { (since } \quad\left[X^{h}, C\right] \in \mathfrak{X}^{v} E\right) .
\end{aligned}
$$

Another characterization of the linear connections is given by the

Proposition 2. $(H C) \Longleftrightarrow \forall X \in \mathfrak{X}(M), \sigma \in \operatorname{Sec} \xi:$

$$
\left(\nabla_{X} \sigma\right)^{v}=\left[X^{h}, \sigma^{v}\right] .
$$

Proof.
(a) Observe first that the mapping

$$
\mathfrak{X}(M) \times \operatorname{Sec} \xi \rightarrow \mathfrak{X}^{v} E, \quad(X, \sigma) \mapsto\left[X^{h}, \sigma^{v}\right]
$$

has the following properties:
(i) it is $\mathbb{R}$-bilinear;
(ii) $\left[(f C)^{h}, \sigma^{v}\right]=(f \circ p)\left[X^{h}, \sigma^{v}\right]$;
(iii) $\left[X^{h},(f \sigma)^{v}\right]=(f \circ p)\left[X^{h}, \sigma^{v}\right]+(X f \circ p) \sigma^{v}\left(f \in C^{\infty}(M)\right)$.

In fact, (i) is obvious since the mappings $X \mapsto X^{h}$ and $\sigma \mapsto \sigma^{v}$ are $\mathbb{R}$-linear, while (ii) and (iii) can be obtained by a straightforward calculation:

$$
\begin{aligned}
{\left[(f X)^{h}, \sigma^{v}\right] } & =\left[(f \circ p) X^{h}, \sigma^{v}\right]=(f \circ p)\left[X^{h}, \sigma^{v}\right]+\sigma^{v}(f \circ p) X^{h}= \\
& =(f \circ p)\left[X^{h}, \sigma^{v}\right]
\end{aligned}
$$

(because $Z \in \mathfrak{X}^{v} E \Longleftrightarrow \forall g \in C^{\infty}(M): Z(g \circ p)=0$ ); the verification of (iii) is similar.
(b) Now we suppose that

$$
\forall X \in \mathfrak{X}(M), \sigma \in \operatorname{Sec} \xi:\left(\nabla_{X} \sigma\right)^{v}=\left[X^{h}, \sigma^{v}\right] .
$$

Applying (ii) we get:

$$
\begin{gathered}
\forall f \in C^{\infty}(M):\left[\nabla_{X}(f \sigma)\right]^{v}=\left[X^{h},(f \sigma)^{v}\right]=(f \circ p)\left[X^{h}, \sigma^{v}\right]+ \\
+(X f \circ p) \sigma^{v}=(f \circ p)\left(\nabla_{X} \sigma\right)^{v}+(X f \circ p) \sigma^{v}=\left[f V_{X}+(X f) \sigma\right]^{v},
\end{gathered}
$$

consequently $\nabla_{X} f \sigma=f \nabla_{X} \sigma+(X f) \sigma$.
It means that the operators $\nabla_{X}$ satisfy the well-known Koszul-axioms of the linear connections, from which (HC) easily follows (see e.g. [16]).
(c) To verify the converse statement, fix a fibered chart for $\xi$, e.g. the chart $\left(p^{-1}(U),\left(x^{i}\right)_{i=1}^{n},\left(y^{\kappa}\right)_{\kappa=1}^{r}\right)$ described in [18], p.1168. Then there exist unique functions

$$
\Gamma_{i}^{\kappa}: p^{-1}(U) \rightarrow \mathbb{R} ; \quad 1 \leq i \leq n, \quad 1 \leq \kappa \leq r
$$

such that

$$
h\left(\frac{\partial}{\partial x^{i}}\right)=\frac{\partial}{\partial x^{i}}-\Gamma_{i}^{\kappa} \frac{\partial}{\partial y^{\kappa}}, \quad h\left(\frac{\partial}{\partial y^{\kappa}}\right)=0 ;
$$

these are the "connection parameters" of $h$. Over $p^{-1}(U)$, we get:

$$
T=\left(y^{\beta} \frac{\partial \Gamma_{i}^{\kappa}}{\partial y^{\beta}}-\Gamma_{1}^{\kappa}\right) d x^{i} \otimes \frac{\partial}{\partial y^{\kappa}}
$$

On the other hand, an easy calculation shows that (locally)

$$
\left[X^{h}, \sigma^{h}\right]=\left(\nabla_{X} \sigma\right)^{v} \Longleftrightarrow \Gamma_{i}^{\kappa}=y^{\beta} \frac{\Gamma_{i}^{\kappa}}{\partial y^{\beta}},
$$

so (HC) implies the condition in question.
Comparing the criterion just obtained with the result of the linearization process sketched in (3.3), we get the following "natural interpretation" of (HC):

Corollary. $h \in \mathcal{H}(\xi)$ satisfies $(H C)$ iff the covariant derivatives belonging to $h$ "essentially" (along sections and up to the canonical mapping $\alpha$ ) coincide with their linearized operators.
(4.3) Suppose that $\nabla \in \mathcal{H}(V \xi)$ is a linear connection, interpreted as a mapping $\mathfrak{X}(E) \times \mathfrak{X}^{v} E \rightarrow \mathfrak{X}^{v} E$ satisfying the Koszul-axioms and let $h \in \mathcal{H}(\xi)$. Then the pair $(\nabla, h)$ is called a Matsumoto-pair (see [17]). The deflection of $(\nabla, h)$ is the vector 1-form

$$
D(\nabla, h): \mathfrak{X}(E) \rightarrow \mathfrak{X}(E), \quad X \mapsto \nabla_{h X} C
$$

where $C$ is the Liouville vector field again. We see immediately that $D(\nabla, h)$ is semibasic. - In particular, choose by way of connection the Berwald connection $\nabla^{B}$ induced by $h$ :

$$
\nabla^{B}: \mathfrak{X}(E) \times \mathfrak{X}^{v} E \rightarrow \mathfrak{X}^{v} E, \quad(X, Y) \mapsto \nabla_{v X}^{i} Y+v[h X, Y]
$$

where $v:=i d-h$ and $\nabla^{i}: \mathfrak{X}^{v} E \times \mathfrak{X}^{v} E \rightarrow \mathfrak{X}^{v} E$ is a pseudoconnection characterized by the following condition:

$$
\forall \sigma \in \operatorname{Sec} \xi, Z \in \mathfrak{X}^{v} E: \nabla_{Z}^{i} \sigma^{v}=0
$$

([18]), Lemma 5). Then $\forall X \in \mathfrak{X}(E)$ :

$$
\begin{aligned}
D\left(\nabla^{B}, h\right)(X) & =\nabla_{h X}^{B} C=v[h X, C]= \\
& =v([h X, C]-h[X, C])=v T(X)=T(X)
\end{aligned}
$$

(since $T$ is also semibasic). It means that in this case deflection and tension coincide. Let us emphasize: deflection plays an important role in Matsumoto's theory of Finsler-connections; see the monograph [9].
(4.4) To conclude our general view on connections, we recall a further basic concept.

Definition 6. The curvature of a connection $h \in \mathcal{H}(\xi)$ is the Nijenhuis torsion $R:=N_{h}=\frac{1}{2}[h, h] . h$ is called flat if $R=0$.

It is easy to check that $R$ is a semibasic 2 -form:
$\operatorname{Im} R \subset \mathfrak{X}^{v} E$ and $R(X, Y)=0, \quad$ if $\quad X \in \mathfrak{X}^{v} E$ or $Y \in \mathfrak{X}^{v} E$.

## 5. The problem of metrizability

(5.1) Now and in the sequel we turn our attention to the tangent bundle case, i.e. we suppose that $\xi:=\tau_{M}=(T M, \tau, M)$ and $h \in \mathcal{H}\left(\tau_{M}\right)$. Then the notion of parallel vector fields along curves has meaning ([7], Def. I.26) and we can raise the following "problem of metrizability":

Find a criterion for the existence of a (smooth) function
(PM) $\quad L: T M \rightarrow \mathbb{R}$ which (in Grifone's terminology) is conserved by parallel transports: for all smooth curve $c: I \rightarrow M$ and parallel vector field $X: I \rightarrow T M, L \circ X: I \rightarrow \mathbb{R}$ is constant.
As we have seen in section 2 , the vector 1-form $h \in \mathcal{H}\left(\tau_{M}\right)$ determines a derivation $d_{h}$ of $\wedge(M)$. The following observation is simple, but very useful:

Lemma 1. ([7]), Prop. I.28) $A$ function $L: T M \rightarrow \mathbb{R}$ is conserved by parallelism if

$$
\begin{equation*}
d_{h} L=0 \tag{6}
\end{equation*}
$$

In coordinates (6) yields the system of partial differential equations

$$
\begin{equation*}
\frac{\partial L}{\partial x^{i}}-\Gamma_{i}^{k} \frac{\partial L}{\partial y^{k}}=0, \quad 1 \leq i \leq n \tag{7}
\end{equation*}
$$

where the $\Gamma_{i}^{k}$-s are the connection parameters of $h$. So our problem is to find the integrability conditions of the system (7). Well now, to attack (PM), we propose the study of formal integrability of the 1st order linear differential operator $d_{h}$ (we shall see soon that $d_{h}$ really has such an interpretation, in a natural manner).
(5.2) By the way, we indicate that (PM) is essentially equivalent with the inverse problem of the calculus of variations. To formulate the latter in modern language, consider a semispray (i.e. a second order differential equation) $S: T M \rightarrow T T M$.

Definition 7. ([8]), p.189) $S$ is said to be variational if there exists a (regular) function $L: T M \rightarrow \mathbb{R}$ such that

$$
i_{S} d d_{J} L=d(L-C L)
$$

where $J: T T M \rightarrow T T M$ is the canonical almost tangent structure (or vertical endomorphism).

By a fundamental result of J. Grifone ([7], Prop. I.41) $h:=\frac{1}{2}(i d+$ $[J, S]$ ) is a connection in $\tau_{M}$. Now it can be shown (see [15]) that
$S$ is variational $\quad \Longleftrightarrow \quad \exists$ (regular) $L: T M \rightarrow \mathbb{R}$ such that $d_{h} L=0, h=\frac{1}{2}(i d+[J, S])$.

From this point of view, an excellent survey and discussion of the problem can be found in the paper [8] of J. Klein.

## 6. The involutivity of $d_{h}$

First we give the promised interpretation of $d_{h}$ as a suitable differential operator.
$—$ Since $\forall X \in \mathfrak{X}(M): d_{h} L(X)=L(h X) \in C^{\infty}(T M), d_{h}$ maps $C^{\infty}(T M)$ into $\operatorname{Sec} \tau_{T M}^{*}$. Here $C^{\infty}(T M)$ can be identified with

$$
\operatorname{Sec} \xi, \quad \xi:=\left(E, p r_{2}, T M\right), \quad E:=\mathbb{R} \times T M
$$

So, in fact, $d_{h}$ can be considered as a 1st order linear differential operator from $\operatorname{Sec} \xi$ into $\operatorname{Sec} \tau_{T M}^{*}$.

## Lemma 2.

(i) The symbol of $d_{h}$ is $s_{0} d_{h}=i_{h}$ (up to an obvious factor which will be omitted).
(ii) At each point $v \in T M$ the null space $\left(g_{0}\right)_{v}:=\operatorname{Ker}\left(s_{0} d_{h}\right)_{v}$ is $n$ dimensional.

Proof. (i) Consider an arbitrary function $f \in C^{\infty}(T M)$ with the condition $f(v)=0$ and a section $\sigma \in \operatorname{Sec} \xi \cong C^{\infty}(T M)$. According to Definition 2,

$$
\begin{aligned}
\left(s_{0} d_{h}\right)(d f(v), \sigma(v)) & =d_{h}(f \sigma)(v)=\left[\left(d_{h} f\right) \sigma+f d_{h} \sigma\right](v)= \\
& =\sigma(v) d_{h} f(v) \stackrel{(2)}{=} \sigma(v) i_{h} d f(v),
\end{aligned}
$$

which gives our assertion.
(ii) Let $\omega \in \mathfrak{X}^{*} T M$ be a 1 -form. By (i),

$$
\omega \in g_{0}=\operatorname{Ker} s_{0} d_{h} \Longleftrightarrow i_{h} \omega=0 \stackrel{(1)}{\Longleftrightarrow} \forall X \in \mathfrak{X}(T M): \omega(h X)=0 .
$$

It means that Ker $s_{0} d_{h}$ is nothing but the annullator of $\mathfrak{X}^{h} T M$, hence $\operatorname{dim} g_{0}=\operatorname{dim} \tau_{T M}-\operatorname{dim} \operatorname{Im} h=2 n-n=n$.

## Lemma 3.

(i) The symbol of the first prolonged operator $p_{1}\left(d_{h}\right)$ is the mapping

$$
s_{1} d_{h}: S^{2} T^{*} T M \otimes E \rightarrow T^{*} T M \otimes T^{*} T M, \quad A \mapsto\left(s_{1} d_{h}\right) A
$$

given by

$$
\left(s_{1} d_{h}\right) A(X, Y)=A(X, h Y) ; \quad X, Y \in \mathfrak{X}(T M)
$$

(ii) $\forall v \in T M: \operatorname{dim}\left(g_{1}\right)_{v}=\frac{n(n+1)}{2}$.

Proof. (i) follows immediately by Definition 2 again, part (i) of the previous Lemma and by (1).
(ii) Let $\left(U,\left(u^{i}\right)\right)$ be a chart for $M$, and let $\left(\tau^{-1}(U) ;\left(x^{i}\right),\left(y^{i}\right)\right)$ be the induced chart for $T M$. Then

$$
\begin{equation*}
\frac{\delta}{\delta x^{i}}:=\frac{\partial}{\partial x^{i}}-\Gamma_{i}^{j} \frac{\partial}{\partial y^{j}}, \quad \frac{\partial}{\partial y^{i}} \quad(1 \leq i \leq n) \tag{8}
\end{equation*}
$$

(the $\Gamma_{i}^{j}$-s are the parameters of $h \in \mathcal{H}\left(\tau_{M}\right)$ ) is a local basis for $\mathfrak{X}(T M)$. Since $g_{1}:=\operatorname{Ker}\left(s_{1} d_{h}\right)$,

$$
A \in g_{1} \Longleftrightarrow \forall X, Y \in \mathfrak{X}(T M): A(X, h Y)=0 .
$$

Using the basis (5), over $\tau^{-1}(U)$

$$
X=X^{i} \frac{\delta}{\delta x^{i}}+X^{n+i} \frac{\partial}{\partial y^{i}}, \quad h Y=Y^{i} \frac{\delta}{\delta x^{i}} ;
$$

so we get:

$$
A \in g_{1} \Longleftrightarrow \begin{cases}A\left(\frac{\delta}{\delta x^{i}}, \frac{\delta}{\delta x^{j}}\right)=0 & (1 \leq i, j \leq n)  \tag{9}\\ A\left(\frac{\delta}{\delta x^{i}}, \frac{\partial}{\partial y^{j}}\right)=0 & (1 \leq i, j \leq n)\end{cases}
$$

Since $A$ is symmetric, (9) gives $\frac{n(n+1)}{2}+n^{2}$ relations for its components, therefore the number of the independent components is

$$
\frac{2 n(2 n+1)}{2}-\frac{n(n+1)}{2}-n^{2}=\frac{n(n+1)}{2}
$$

Clearly, this is just the dimension of $g_{1}$.
Proposition 3. The operator $d_{h}$ is involutive.
Proof. Let $\left(\tau^{-1}(U) ;\left(x^{i}\right),\left(y^{i}\right)\right)$ be an induced chart for $T M$. First we note that

$$
\delta y^{i}:=\Gamma_{j}^{i} d x^{j}+d y^{i}, \quad d x^{i} \quad(1 \leq i \leq n)
$$

is the dual of the local basis (8), so (over $\tau^{-1}(U)$ )

$$
\forall \omega \in \mathfrak{X}^{*} T M: \omega=\omega_{i} \delta y^{i}+\omega_{n+i} d x^{i} .
$$

In particular, if $\omega \in g_{0}$, then (as we have seen) $\omega$ annihilates $\mathfrak{X}^{h} T M$, hence

$$
\forall 1 \leq i \leq n: \quad 0=\omega\left(\frac{\delta}{\delta x^{i}}\right)=\omega_{n+j} d x^{j}\left(\frac{\delta}{\delta x^{i}}\right)=\omega_{n+i}
$$

i.e.

$$
\omega \in g_{0} \Longleftrightarrow \omega=\omega_{i} \delta y^{i}
$$

Now we are going to show that $\forall v \in \tau^{-1}(U)$ :

$$
\left(\left(\frac{\partial}{\partial y^{i}}\right)_{v},\left(\frac{\delta}{\delta x^{i}}\right)_{v}\right) \quad 1 \leq i \leq n
$$

is a quasi-regular basis. In fact, using the preceding remark,

$$
\left(g_{0}\right)_{\frac{\partial}{\partial y^{1}}} \stackrel{\text { Def. } 3}{=}\left\{\omega \in g_{0} \left\lvert\, i_{\frac{\partial}{\partial y^{1}}} \omega=0\right.\right\}=\left\{\omega=\omega_{i} \delta y^{i} \in g_{0} \mid \omega_{1}=0\right\}
$$

which means that

$$
\operatorname{dim}\left(g_{0}\right) \frac{\partial}{\partial y^{\mathrm{I}}}=n-1
$$

Similarly,

$$
\begin{gathered}
\left(g_{0}\right)_{\left(\frac{\partial}{\partial y^{1}}, \frac{\partial}{\partial y^{2}}\right)}=\left\{\omega=\omega_{i} y^{i} \in g_{0} \mid \omega_{1}=\omega_{2}=0\right\} \Rightarrow \\
\operatorname{dim}\left(g_{0}\right)_{\left(\frac{\partial}{\partial y^{1}}, \frac{\partial}{\partial y^{2}}\right)}=n-2 .
\end{gathered}
$$

Proceeding in the same way, finally we get:

$$
\operatorname{dim}\left(g_{0}\right)_{\left(\frac{\partial}{\partial y^{1}}, \ldots, \frac{\partial}{\partial y^{n-1}}\right)}=1
$$

By these

$$
\sum_{j=1}^{n-1} \operatorname{dim}\left(g_{0}\right)_{\left(\frac{\partial}{\partial y^{1}}, \ldots, \frac{\partial}{\partial y^{j}}\right)}=\sum_{j=1}^{n-1}(n-j)=\frac{n(n-1)}{2}
$$

consequently

$$
\begin{gathered}
\operatorname{dim}\left(g_{0}\right)+\sum_{j=1}^{2 n} \operatorname{dim}\left(g_{0}\right)_{\left(e_{1}, \ldots, e_{j}\right)} \stackrel{\text { Lemma } 2}{=} n+\frac{n(n-1)}{2}= \\
=\frac{n(n+1)}{2} \stackrel{\text { Lemma }^{3}}{=} \operatorname{dim}\left(g_{1}\right) \\
\left(e_{1}:=\frac{\partial}{\partial y^{i}}, e_{n+i}:=\frac{\delta}{\delta x^{i}} ; \quad 1 \leq i \leq n\right)
\end{gathered}
$$

which means the involutivity of $d_{h}$.

## 7. The surjectivity of $\pi_{2}$

Keeping the previous notations and conventions, consider the following diagram:


Here $K:=\frac{T^{*} T M \otimes T^{*} T M}{\operatorname{Im} s_{1} d_{h}}$ is said to be the space of obstructions. By the "algorithm" sketched in [2], in order to verify the surjectivity of $\pi_{2}: R_{2} \rightarrow R_{1}$ one has to
(i) give a "good interpretation" for the space of obstructions and construct a mapping $\tau$ which makes the diagram commutative;
(ii) define a linear connection $\nabla$ in $\tau_{T M}^{*}$ such that

$$
\forall v \in T M, \sigma \in \operatorname{Sec} \xi: d_{h} \sigma(v)=0 \Rightarrow \tau\left(\nabla d_{h} \sigma\right)_{v}=0
$$

(To realize (ii), it is sufficient to give a suitable connection in $\tau_{T M}$ since it induces the desired connection in $\tau_{T M}^{*}$.)
Of course, in our case the surjectivity of $\pi_{2}$ does not hold without further assumptions. In what follows, we are going to prove this only under the quite restrictive condition $N_{h}=0$.

Lemma 4. Let $v:=i d-h$, as above. The mapping $v: T T M \rightarrow$ $T T M$ induces a pull-back map $v^{*}: \wedge^{2} T M \rightarrow \wedge^{2} T M$. - The space of obstructions is (pointwise) isomorphic to the space

$$
\operatorname{Ker} v^{*} \oplus \tau_{2}^{0}\left(T^{v} T M\right)
$$

$\left(\tau_{2}^{0}\left(T^{v} T M\right)\right.$ is the space of $(0,2)$ tensors $\left.\mathfrak{X}^{v} T M \times \mathfrak{X}^{v} T M \rightarrow C^{\infty}(T M)\right)$.
Proof. On the one hand, we have:

$$
\begin{aligned}
\operatorname{dim} K & =\operatorname{dim}\left(T^{*} T M \otimes T^{*} T M\right)-\operatorname{rank} s_{1} d_{h}= \\
& =(2 n)^{2}-\left(\operatorname{dim} S^{2} T^{*} T M-\operatorname{dim} g_{1}\right) \stackrel{\text { Lemma } 3}{=} \\
& =4 n^{2}-\frac{2 n(2 n+1)}{2}+\frac{n(n+1)}{2}=\left(n^{2}+\frac{n(n-1)}{2}\right)+n^{2} .
\end{aligned}
$$

On the other hand,

$$
\begin{aligned}
\operatorname{dim} \operatorname{Ker} v^{*}=\operatorname{dim} \wedge^{2} T M & -\operatorname{dim} \operatorname{Im} v^{*}=\frac{2 n(2 n-1)}{2}-\frac{n(n-1)}{2}= \\
& =n^{2}+\frac{n(n-1)}{2}
\end{aligned}
$$

while $\operatorname{dim} \tau_{2}^{0}\left(T^{v} T M\right)=n^{2}$, from which the assertion follows.
Proposition 4. Let the mappings

$$
\begin{aligned}
& \tau_{1}: T^{*} T M \otimes T^{*} T M \rightarrow \operatorname{Ker} v^{*} \quad \text { and } \\
& \tau_{2}: T^{*} T M \otimes T^{*} T M \rightarrow \tau_{2}^{0}\left(T^{v} T M\right)
\end{aligned}
$$

be defined as follows:

$$
\begin{gathered}
\forall B \in T^{*} T M \otimes T^{*} T M: \\
\tau_{1} B(X, Y):=B(h X, Y)-B(h Y, X) \quad(X, Y \in \mathfrak{X}(T M)), \\
\tau_{2} B(U, V):=B(U, V) \quad\left(U, V \in \mathfrak{X}^{v} T M\right) .
\end{gathered}
$$

Then the sequence

$$
S^{2} T^{*} T M \xrightarrow{s_{1} d_{h}} T^{*} T M \otimes T^{*} T M \xrightarrow{\tau:=\tau_{1} \oplus \tau_{2}} \operatorname{Ker} v^{*} \oplus \mathcal{T}_{2}^{0}\left(T^{v} T M\right) \rightarrow 0
$$

is exact
Proof. (a) We show that $\operatorname{Im} s_{1} d_{h} \subset \operatorname{Ker} \tau$. In fact,

$$
\begin{gathered}
\forall B \in S^{2} T^{*} T M ; \quad \forall X, Y \in \mathfrak{X}(T M): \\
{\left[\left(\tau_{1} \circ s_{1} d_{h}\right) B\right](X, Y)=s_{1} d_{h}(B)(h X, Y)-s_{1} d_{h}(B)(h Y, X)=} \\
\text { Lemma }{ }^{3,(\mathrm{i})} B(h X, h Y)-B(h Y, h X)=0 \\
{\left[\left(\tau_{2} \circ s_{1} d_{h}\right) B\right](v X, v Y)=s_{1} d_{h}(B)(v X, v Y)=B(v X, h v Y)=0 .}
\end{gathered}
$$

(b) $A \in \operatorname{Ker} \tau \Longleftrightarrow A \in \operatorname{Ker} \tau_{1} \wedge A \in \operatorname{Ker} \tau_{2}$. Condition $A \in \operatorname{Ker} \tau_{1}$ means that

$$
A\left(\frac{\delta}{\delta x^{i}}, \frac{\delta}{\delta x^{j}}\right)=A\left(\frac{\delta}{\delta x^{j}}, \frac{\delta}{\delta x^{i}}\right) \wedge A\left(\frac{\delta}{\delta x^{i}}, \frac{\partial}{\partial y^{j}}\right)=0 \quad(1 \leq i, j \leq n)
$$

while $A \in \operatorname{Ker} \tau_{2} \Longleftrightarrow A\left(\frac{\partial}{\partial y^{i}}, \frac{\partial}{\partial y^{j}}\right)=0 \quad(1 \leq i, j \leq n)$. Hence the elements of $\operatorname{Ker} \tau$ can be characterized by $n^{2}+\frac{n(n+1)}{2}$ independent components with respect to the local basis (8), consequently $\operatorname{dim} \operatorname{Ker} \tau=$ $n^{2}+\frac{n(n+1)}{2}$. But
$\operatorname{rank} s_{1} d_{h}=\operatorname{dim} S^{2} T^{*} T M-\operatorname{dim} \operatorname{Ker} s_{1} d_{h} \stackrel{\text { Lemma }}{=}$ 3, (ii) $n^{2}+\frac{n(n+1)}{2}$.

Combining these with (a), we get: $\operatorname{Im} s_{1} d_{h}=\operatorname{Ker} \tau$.
(c) Finally, since

$$
\begin{aligned}
\operatorname{rank} \tau & =\operatorname{dim}\left(T^{*} T M \otimes T^{*} T M\right)-\operatorname{dim} \operatorname{Ker} \tau \stackrel{(\mathrm{b})}{=}(2 n)^{2}-\left(n^{2}+\frac{n(n+1)}{2}\right)= \\
& =2 n^{2}+\frac{n(n-1)}{2} \stackrel{\text { Lemma }^{4}}{=} \operatorname{dim} K=\operatorname{dim} \operatorname{Ker} V+\operatorname{dim} \mathcal{T}_{2}^{0} T^{v} T M
\end{aligned}
$$

$\tau$ is also surjective, which concludes the proof.
Lemma 5. Let $\nabla$ be a linear connection in $\tau_{T M}$ with torsion $T$. Then for each 1-form $\omega \in \wedge^{1} T M$, we have:

$$
\begin{aligned}
\tau_{1}(\nabla \omega)(X, Y)= & d_{h} \omega(X, Y)-\omega\left(\nabla_{Y} h X-\nabla_{X} h Y+T(h X, Y)+\right. \\
& +T(X, h Y)+h[X, Y])
\end{aligned}
$$

Proof. Using the rules of Frölicher-Nijenhuis calculus, the covariant differentiation and the definition of $T$, we get:

$$
\begin{aligned}
d_{h} \omega(X, Y)= & {\left[\left(i_{h} \circ d-d \circ i_{h}\right) \omega\right](X, Y)=\left[\left(i_{h} d \omega\right)(X, Y)-d\left(i_{h} \omega\right)\right](X, Y)=} \\
= & d \omega(h X, Y)+d \omega(X, h Y)-X\left(i_{h} \omega\right)(Y)+ \\
& +Y\left(i_{h} \omega\right)(X)+\omega(h[X, Y])= \\
= & h X \omega(Y)-h Y \omega(X)-\omega([h X, Y]+[X, h Y]-h[X, Y])= \\
= & h X \omega(Y)-h Y \omega(X)-\omega\left(\nabla_{h X} Y-\nabla_{Y} h X-T(h X, Y)+\right. \\
& \left.+\nabla_{X} h X-\nabla_{h Y} X-T(X, h Y)-h[X, Y]\right)= \\
= & \nabla \omega(h X, Y)-\nabla^{2} \omega(h Y, X)+ \\
& +\omega\left(\nabla_{Y} h X-\nabla_{X} h Y+T(h X, Y)+T(X, h Y)+h[X, Y]\right) .
\end{aligned}
$$

Since $\nabla \omega(h X, Y)-\nabla \omega(h Y, X)=\left[\tau_{1}(\nabla \omega)\right](X, Y)$, we have the desired formula.

Corollary. If $\omega_{v}=0$, then $\forall v \in T M: \tau_{1}(\nabla \omega)_{v}=\left(d_{h} \omega\right)_{v}$.
Lemma 6. Let $\nabla$ be a linear connection in $\tau_{T M}$. If $f \in C^{\infty}(T M)$ $(\cong \operatorname{Sec} \xi)$ and $\left(d_{h} f\right)_{x}=0(x \in T M)$, then $\tau_{2}\left(\nabla d_{h} f\right)_{x}=0$.

Proof. $\forall U, V \in \mathfrak{X}^{V} T M: \tau_{2}\left(\nabla d_{h} f\right)(U, V):=\nabla d_{h} f(U, V)=$

$$
=\left(\nabla_{U} d_{h} f\right)(V)=U\left[d_{h} f(V)\right]-d_{h} f\left(\nabla_{U} V\right)
$$

Since $\left(d_{h} f\right)_{x}=0$, it follows that

$$
\left[\tau_{2}\left(d_{h} f\right)\right]_{x}\left(U_{x}, V_{x}\right)=\left[U\left(d_{h} f(V)\right)\right]_{x} \stackrel{(2)}{=} U_{x}\left[i_{h} d f(V)\right]=U_{x}[d f(h V)]
$$

But $V \in \mathfrak{X}^{v} T M \Rightarrow h V=0$, so $\tau_{2}\left(d_{h} f\right)_{x}=0$.
Piecing together our previous findings, we arrive at the main result of the paper.

Theorem. If $h \in \mathcal{H}\left(\tau_{T M}\right)$ is a flat (nonlinear) connection, then the differential operator $d_{h}$ is formally integrable.

Proof. We have already known that $d_{h}$ is involutive (Proposition 3). If $\nabla$ is a linear connection in $\tau_{T M}$ and $\left(d_{h} f\right)_{x}=0$, then in view of the Corollary of Lemma 5,

$$
\tau_{1}\left(d_{h} f\right)_{x}=\left(d_{h} d_{h} f\right)_{x}=\left(d_{h}^{2} f\right)_{x} \stackrel{(5)}{=} 0
$$

since $[h, h]=0$ by the flatness. Combining this with Lemma 6, we get that $\tau\left(d_{h} f\right)_{x}=0$. By our above remarks this guarantees that condition $2^{\circ}$ of the Cartan-Kähler-Goldschmidt theorem is also satisfied.

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J. SZILASI and Z. MUZSNAY

MATHEMATICAL INSTITUTE
OF KOSSUTH LAJOS UNIVERSITY
DEBRECEN, H-4010
P.O. BOX 12
(Received September 30, 1992)

